Towards Scale-invariant Cyclic universes

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arXiv:0801.1315 [hep-th] arXiv:0707.4679 [hep-th] with S Alexander & R Brandenberger

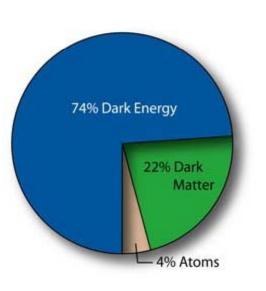
'Hawk 1998

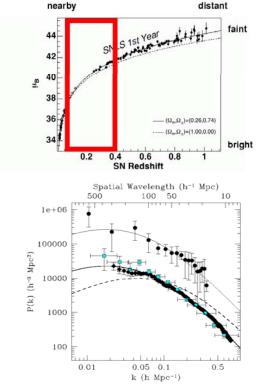
Standard Model of Cosmology

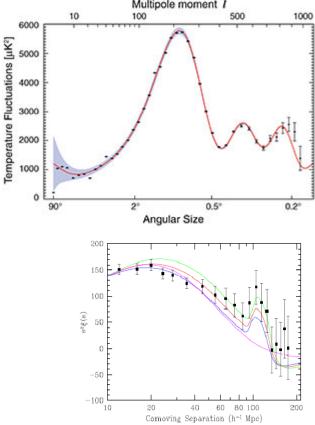
$$ds^{2} = a^{2}(\tau)[-(1 + \Phi(\tau, x))d\tau^{2} + (1 - \Phi(\tau, x))dx^{2}]$$

Key ingredients $k^{3}P_{\Phi}(k) \approx const$

- > GR+Homogeneous Isotropic cosmology+ Perturbations
- > Inflation → Radiation → Matter(DM+baryons)→Dark Energy Concordant ∧CDM model







Why do we still have our jobs?

Finally we know for sure that we almost know nothing Inflaton, Baryons, Dark Matter, Dark Energy????

Two Attitudes

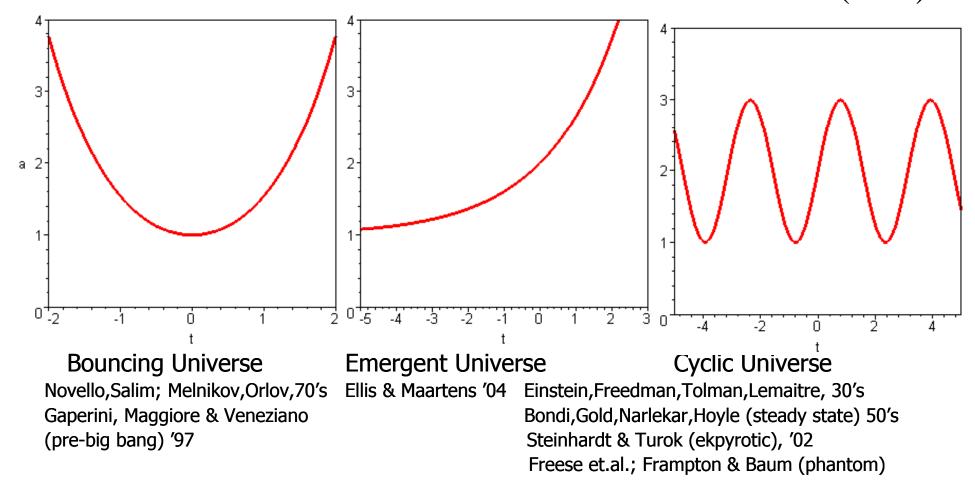
- > SM is fine, try to address the hard questions
- Search for alternative Non-singular (QG) cosmology
- Cosmology is an ultimate magnifier
- 1. Length

$$H_0^{-1} \xrightarrow{\text{radiation}} \frac{T_0}{T_{GUT}} H_0^{-1} \xrightarrow{\text{inflation}} e^{-N} \left(\frac{M_p^2}{T_{GUT}T_0}\right) l_p = e^{-(N-75)} l_p$$

- 2. Energy density can become Plackian
- 3. Spatial curvature can become Planckian
- Observable Quantum Gravity Effects may be testable already (WMAP) and/or near future (Planck)
- 1. Spectrum
- 2. Gravity waves
- 3. Non gaussianity and non adiabaticity

Non-singular Cosmologies

- "Effective" FLRW cosmologies a good description: LQC + BKL \succ $\vec{R} \sim H^2 \sim (\dot{a}/a)^2$
- Look into the "eternal past" \succ



Plan of the talk

- > Virtues and Challenges of cyclic cosmology
- > Tracking perturbations around Bounce
- > New "emergent cyclic Universe"
- > Cyclic Inflation

<u>Challenges & Virtues of Cyclic</u> <u>Universes</u>

I. Nonsingular geodesically complete Universe

- Consistent (ghostfree) Bounce
- LQC
- p-adic/SFT inspired Non-local modifications of gravity [Seigel, Mazumdar, YT]
- BCS Gap energy: [Alexander, & YT]
- Coupling of fermions to gravity \Rightarrow four fermion interaction
- Attractive \Rightarrow negative energy required for bounce
- Gap energy \rightarrow chemical potential \rightarrow number density \rightarrow volume
- Nontrivial volume dependence can temporarily violate DEC

II. Turn-around

- Spatial curvature (no Dark Energy)
- > Scalars, phantom

III. Classic puzzles:

- > Horizon: All
- > Flatness, Largeness, Entropy: Emergent cyclic models
- > 7 Homogeneity/Isotropy: Some aspects we will be able to address

IV. Black hole over-production & Dark energy

- > Quintessence (ekpyrotic) makes BH's dilute
- Phantom makes BH's disintegrate

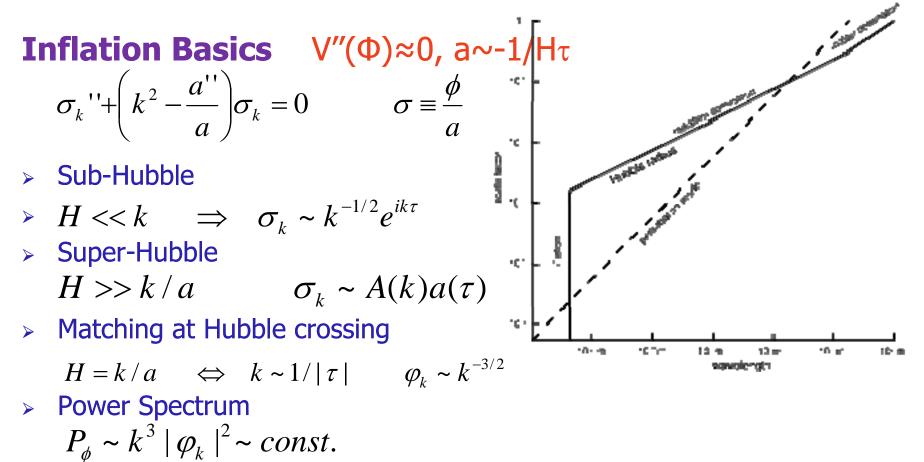
v. ? Generating Scale-invariant Perturbations

- > Scalar field fluctuations [Steinhardt et.al]
- Stringy thermal fluctuations [Brandenberger et. al]
- > "New" ideas

vi. Tolman's Entropy or the "Super Big Bang" Problem

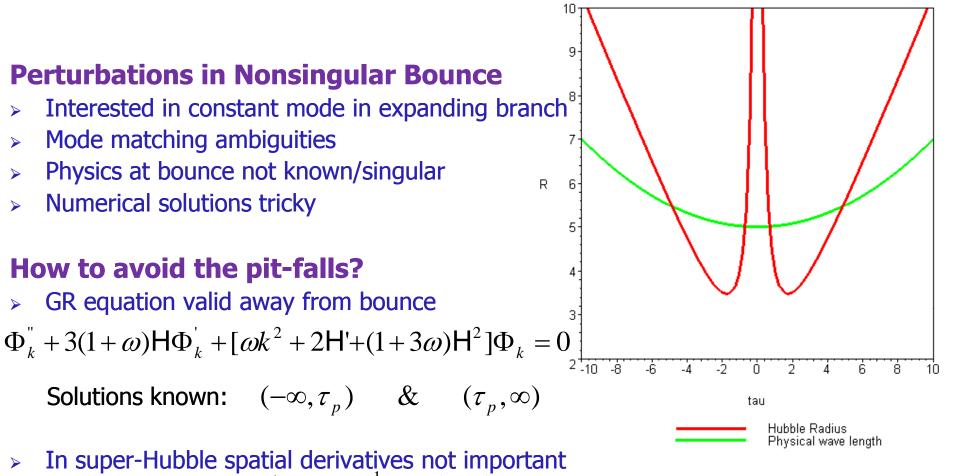
- > Thermal UV phase (Hagedorn phase) exists +
- BB singularity problem is solved

Transferring fluctuations via bounce



Ekpyrotic

 Generates scale-invariant spectrum in the "growing mode" during contraction phase.



- In super-Hubble spatial derivatives not important If no inflation $k_{ph} \ll \tau_p^{-1} \sim O(M_p)$ If no inflation .
- Near bounce also like super-Hubble , even if $H \rightarrow 0$ •
- Higher order spatial derivatives can also be neglected
- Perturbation eqn. only depends on time,
- find fluctuations for scale-factor \Leftrightarrow only need homogeneous isotropic cosmology
- Matching in overlapping regions \succ $(-\tau_k, -\tau_n)$ & (τ_n, τ_k)
- $H' \ll \tau_n^{-1}$

 $(-\tau_k, \tau_k)$

An example

Assumption

- > Perturbation equation not effected by new physics
- > Non-local physics changes global evolution
- > Does not effect perturbations
- Casimir Energy, Gap energy

Results

- Depends on the new physics (background)
- Generally modes mix! Good news for ekpyrotic scenarios
- Modes can even switch
- Modes switch at $\omega \approx 5.314$,
- Mixes when $\omega >>1$ (ekprotic limit)

Emergent Cyclic Universe

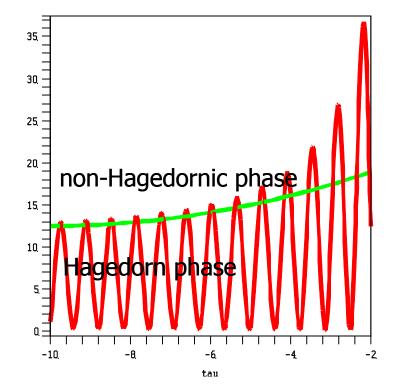
Tolman's Entropy Problem

- Entropy is monotonically increasing
- Universe atmost quasi-periodic
- Entropy (Energy, period) vanishes in a finite time in the past
- Beginning of time back to square 1
- > Thermal Hagedorn Phase, $T=T_H \sim M_S$
- All string states in thermal equilibrium entropy constant
- Below critical temperature, massive states decouple, entropy produced.
- As cycles shrink, universe is hotter less time in entropy producing

more time in

 $\Delta S \to 0 \qquad \Longrightarrow S_n \to S_{-\infty} \neq 0$

Cycles asymptote to a periodic evolution





Stringy Toy Model

> Hagedorn half
$$H^2 = \frac{1}{3} [\rho_{hag} - \rho_{cas} + \rho_{curv}] = \frac{1}{3} \left[\frac{S}{a^3} - \frac{\Omega_c}{a^4} - \frac{\Omega_k}{a^2} \right]$$

• Eternally periodic universe

$$a(\tau) = \frac{S - \sqrt{S^2 - 4\Omega_c \Omega_k \cos v\tau}}{2\Omega_k} \qquad v \equiv \frac{T_H^2}{M_p} \sqrt{\frac{\Omega_k}{3}}$$

- Non-Hagedorn half (no energy exchange)
- Hagedorn matter = massless (r)+massive (m)
- $S \to \Omega_m \quad \& \quad \Omega_c \to \Omega_c \Omega_r \quad \Rightarrow \quad a_{\max} \sim \sqrt{\Omega_r}$
- > Gluing the two halves (Ω_r , Ω_m , a_{tran})
- Entropy is conserved
- Matter and radiation was in equilibrium till the transition

• Phenomenological Input:
$$\mu \equiv \frac{\rho_m}{\rho_r} \sim 10^{-22} << 1$$

$$\Omega_r \sim S^{4/3} \Longrightarrow a_{\max} \sim S^{2/3} \qquad a_{tran} \sim S^{1/3}$$

Energy Exchange & Entropy Production

 T_{H}

 $S = S_{cr} \left| 1 + \frac{1}{Cn^2} \right|$

Matter coverted to radiation [Tolman, Barrow et.al.] $\dot{\rho}_r + 4H\rho_r = T_r^4 s$ $\dot{\rho}_m + 4H\rho_m = -T_H^4 s$

- Consistent with 1st & 2nd
- **Conserves total**
- Breaks time-reversal symmetry
- Exchange function, $s(a, \Omega)$
- Small cycle limit, s ~ const.

law of thermodynamics stress-energy tensor arrow of time

Crucial Difference

- Singular bounce 📄 Г interaction canno

keep up with $H \rightarrow \infty$ Thermal equil cannot be mantiened is entropy production during bounce Nonsingular bounce 📄 H is finite, a thermal Hagedorn phase can exist.

Cyclic Inflation

Can we mimick inflation?

Entropy production by decay of massive particles:

$$\left(\frac{S_{n+1}}{S_n}\right) = \left(\frac{S_r}{S_m}\right) \sim \frac{\rho_r^{3/4}V}{\rho_m V M^{-1}} \sim \frac{T_H}{T_d} \equiv \kappa$$

- Entropy increases by a constant factor
- If decay time > Period, κ smaller
- So does the scale factor!
- Energy density remains same

$$a_{tran} \sim S^{1/3} \Rightarrow \frac{a_{n+1}}{a_n} \sim \kappa^{1/3}$$

Inflation over many many cycles

> Massless scalar field will see inflation on an average, if

$$\tau H_{av} \ll 1$$

More detail requirements

- With curvature turnaround τ increases
- $\begin{array}{l} & \text{-ve consmological constant} \\ & H_{av} = \frac{\int H dt}{\int dt} = \frac{\ln\left(\frac{a_{n+1}}{a_n}\right)}{\tau} \Longrightarrow \ln\left(\frac{a_{n+1}}{a_n}\right) = \frac{\ln \kappa}{3} <<1 \end{array}$
- Over many many many cycles we realize inflation \succ

How to exit?

- Introduce a potential
- Kinetic energy at end of contraction depends on slope:
- Total energy can go from ve to + ve
- Can zoom past the minimum and into + ve potential region
- Can even mimick quintessence, not necessary.

Thermal fluctuations

Spectrum [Peebles'93, Pogosian & Magueijo]
Mechanism
$$\delta_L^2 = \frac{\Delta E^2}{E^2}\Big|_L = \left(\frac{d^2 \ln Z}{d\beta^2}\right) / \left(\frac{d \ln Z}{d\beta}\right)^2 = \frac{T^2}{\rho^2 L^3} \frac{\partial \rho}{\partial T} \sim \frac{1}{(TL)^3}$$

Energy fluctuation in a given volume related to Power spectrum

• White noise spectrum
$$\delta_L^2 = \int dk k^2 W(kL) P(k) \Rightarrow P(k) = \delta_k^2 = \frac{T^2}{\rho^2} \frac{\partial \rho}{\partial T} k^0 \rightarrow \frac{4}{g} T^{-3}$$

- Thermal correlations only for sub-Hubble modes $\frac{k}{a} \sim H \sim \frac{T^2}{M_n} \Longrightarrow k \sim T \Longrightarrow P(k) = k^{-3}$
- Hubble Crossing
- Conformal radiation leads to scale invariance
- P(k) transferred to gravity, $P_{\Phi}(k)$ $\nabla^2 \Phi \approx \delta \rho \Longrightarrow k^2 \Phi_{\mu} = a^2 \rho \delta_{\mu} \sim a^2 H^2 \delta_{\mu}$
- Super Hubble governed by Φ_k
- Robust and general mechanism (no fine-tuning): \geq
- At Hubble crossing validity of GR
- Radiation domination
- Contraction so that modes leave the Hubble radius
- No mode mixing, Φ_k remains constant

Amplitude

- > The problem for symmetric bounce
- > Consider scale which enters Hubble radius at T_{eq} (1Mpc)

$$\delta_L^2 = \frac{1}{g_* (LT_L)^3} \sim \frac{1}{g_*} \left(\frac{H_L}{T_L}\right)^3 \sim \sqrt{g_*} \left(\frac{T_L}{M_p}\right)^3 <<10^{-8} \rightarrow T_L \sim 10^{13} Gev$$

Exactly the right scenario to "amplify" perturbations

$$k = T_L^2 a_L = T_{eq}^2 a_{eq} \qquad \& \qquad S \sim (Ta)^3 \sim T^{-3}$$
$$\left(\frac{S_0}{S_{-1}}\right) = \kappa = 10^{67}$$

- a_p is increasing λ_{ph} is decreasing with cycles \Rightarrow universe has to expand more for the same mode to exit we need

- > Thermal Equilibrium in Hagedorn phase [Frey et.al, Hindmarsh & Skliros]
- Small cycles: Thermal equilibrium to be mantained

$$\Gamma_{\rm int} \sim T_H g_s^2 n \qquad H \sim T_H^2 / M_p$$

 $T_H \sim 10^{-3} M_p \qquad g_s^2 > 10^{-5}$

• Large cycles:
$$H \propto S^2$$

Thermal equilibrium around bounce short lived Near bounce large amounts of entropy produced

- > Entropy Production around Bounce
- Naïve Estimate: when radiation converts to Hagedorn matter

$$\left(\frac{S_0}{S_{-1}}\right) = \left(\frac{S_H}{S_r}\right) \sim \frac{\rho_H V T_H^{-1}}{\rho_r^{3/4} V} \sim \frac{\rho_b^{1/4}}{T_H} \sim S_{-1}$$

 Holographic Saturation: Can't trust physics at super-Planckian energy densities Treat bounce at black box, use holographic entropy bound If saturated we get the same estimate!

$$S_0 \sim 10^{134}$$
 & $S_{-1} \sim 10^{67}$

$$S \leq area \sim \left(\frac{E}{M_p}\right)^2 \Longrightarrow S_0 \leq S_{-1}^2 \left(\frac{T_H}{M_p}\right)^2$$

 $H < \Gamma_{\rm int}$

Dark Energy and Black Hole Problem

- > Add $\Lambda \sim$ (mev) there exists a last (our) cycle! DE phase
- > Previous cycles very short
- no matter domination, no LSS
- no junk (BH/inhomogeneities) from previous cycle
- thermal density fluctuation very small
- > Toy Model can be extended to ekpyrotic/phantom type late cycles

Summary

The story so far?

- Began as string size, curved universe, in almost periodic
 & almost Hagedornic phase: emerging phase
- > (a) Constant cycle "inflationary" phase -> graceful exit to long lived cycle
- > (b) With entropy production, Hagedorn phase becomes shorter,
- large entropy production starts
- Universe highly asymmetric and large
- Very large cycles (like ours), curvature gives way to Dark Energy

Problems addressed

- > Singularities: BBS & super BBP
- > Classic Problems: horizon, flatness, entropy, largeness
- > Late time: inhomogeneity/BH overproduction & DE
- > Perturbations: spectrum and amplitude

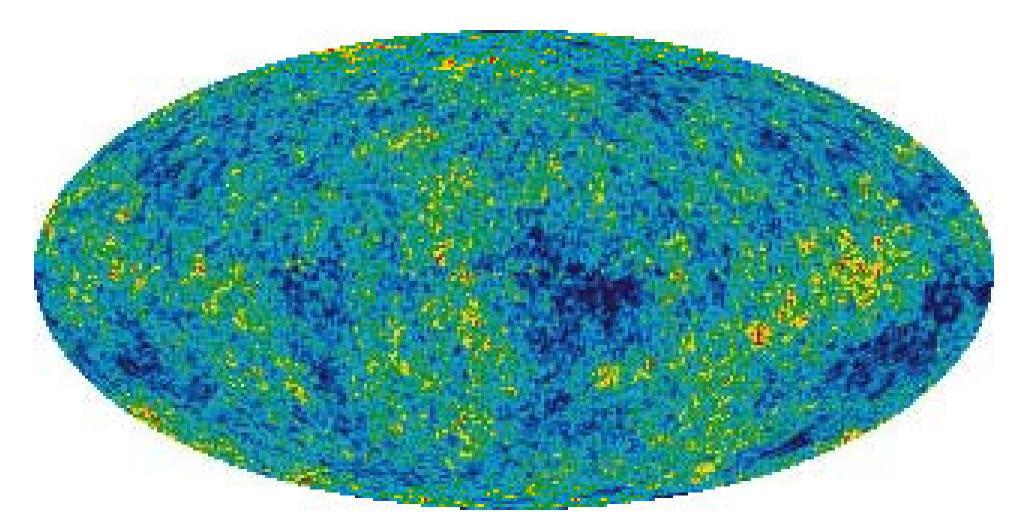
Needs investigation

- > Initial homogeneity/isotropy, Mix master behaviour
- > Hagedorn physics: Non-equilibrium dynamics, Casimir Energy, Decay rate...
- > Mode mixing during Bounce? Is Φ really constant during bounce?

Predictions/Tests

> Gravity Waves, Non-gaussianity...

Precision Cosmology



Cosmic Microwave Background

We need an asymmetric bounce, the comoving scale has to exit the Hubble radius much earlier.

 This is precisely what we have via entropy production

Quantum Gravity: A toy Model

Motivation

Stringy

Dual Field theory action for strings on Random Lattice

$$\hat{S} = \int d^D x \ tr \left[\frac{1}{2} \phi e^{-\alpha' \Box/2} \phi + G^{n-2} \phi^n \right]$$

Linear Regge trajectories: Confinement [Grisaru, Siegel, Y.T.]

- Tachyons in open SFT and p-adic string theory has similar form
- Higher Derivative but Ghost free

$$S = \int d^4x \, \phi \Box (\Box + m^2) \phi \Rightarrow \Box (\Box + m^2) \phi = 0$$

$$\Delta(p^2) = \frac{1}{p^2(p^2 + m^2)} \sim \frac{1}{p^2} - \frac{1}{(p^2 + m^2)}$$

- Non-singular UV & IR behavior Weinberg's "Asymptotic safety"
- Non-perturbative Quantum gravity
 Close to Planck scale, all terms important

Model Action $S = \frac{M_p^2}{2} \int d^4x \sqrt{-g} \left[R + \sum_{n=0}^{\infty} c_n R \Diamond^n R \right] \qquad c_n \sim \frac{1}{M^{2n+2}}$

Generalized ``Einstein's" Field Equations

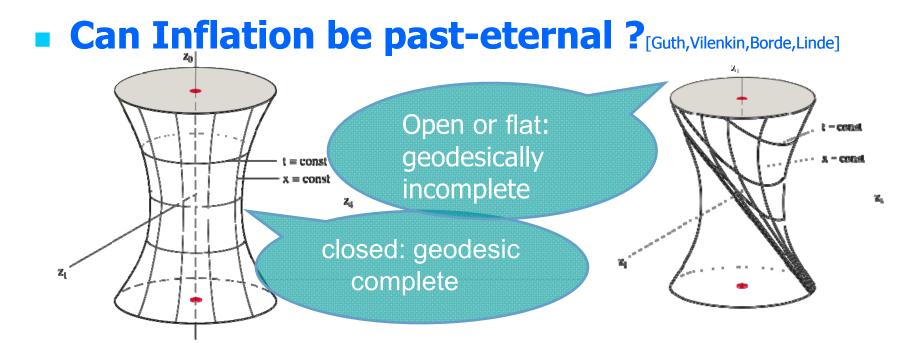
$$\bar{G}_{\mu\nu} \equiv G_{\mu\nu} + \sum_{n=0}^{\infty} G_{\mu\nu}^n = T_{\mu\nu}$$

$$G^n_{\mu\nu} \sim c_n(\Box^{n+1}R + \Box^p R \Box^m R)$$

• Conservation Equation $\nabla^{\mu} \tilde{G}_{\mu\nu} = 0$

For Cosmology $\tilde{G}_{00} = 0$ equation suffices Exact Bouncing solutions! $a(t) = \cosh(\lambda t)$

$$\Lambda \neq 0, \qquad \rho_{rad} \sim M_s^2 M_p^2 \qquad \& \quad \lambda \sim M_s$$



• **Big Bounce:** H = 0 & $\ddot{a} > 0$ $H^2 = \frac{\sum \rho_I}{3M_p^2}$

Flat/Open: DEC,WEC violation

<u>Plan for the rest of the Talk</u>

- Consistent (ghostfree) modification of Gravity LQC, Bouncing and cyclic Universes Non-perturbative gravity ⇒ Bouncing Universe
- Consistent (ghostfree) modification of Matter-sector
- Ghost Condensation [Arkani-Hamed et.al., Khoury et.al.]
- Casimir Energy and an "emergent cyclic Universe Hagedorn Physics & Tolman's Entropy Problem
- Gap Energy and cosmological BCS Condensation
- Generating Scale-invariant Perturbations
- Scalar field fluctuations [Steinhardt et.al]
- Stringy thermodynamic fluctuations [Brndenberger,Nayeri & Vafa]

[Alexander & Vaid]

Thermodynamic Fluctuations during Hagedorn Phase

Energy to Power Spectrum

Energy to density fluctuations:

$$|\delta \rho_k|^2 \sim k^3 |\delta E(r \sim k^{-1})|^2$$

Energy fluctuations from Heat Capacity

$$Z \sim \sum e^{-\beta E} \qquad \Rightarrow \delta E(r) \sim T^2 C_v \sim \frac{T}{T_H - T} r^2$$

Matter fluctuations to metric perturbations

$$\nabla^2 \Phi = 4\pi G \delta \rho \qquad \Rightarrow |\Phi_k|^2 \sim k^{-4} |\delta \rho_k|^2$$

Power Spectrum

$$P_{\Phi} \sim k^3 |\Phi_k|^2 \sim \frac{T}{T_H - T}$$

CMB Spectrum: Minimal Requirements

 Fluctuations should come from massive modes and not massless (radiation)

• Amplitude: $C_{rad} << C_{massive} \qquad \Rightarrow \frac{\Delta T}{T_{H}} < 10^{-30}$ $\delta_{CMB}^{2} \sim 10^{-10} \sim \left(\frac{M_{s}}{M_{p}}\right)^{4} \frac{T_{H}}{\Delta T}$ $\frac{T_{H}}{\Delta T} = 10^{30} \qquad \Rightarrow \frac{M_{s}}{M_{p}} \sim 10^{-10}$ • Spectral tilt: $|\eta_{s} - 1| \approx 10^{-60} \left(\frac{M_{s}}{\lambda}\right)^{2} \frac{T_{H}}{\Delta T}$

typically one obtains almost perfect scale-invariance!

- t' Hooft dual to string theory
- Polyakov action: $S = \int \frac{d^2\sigma}{2\pi} \sqrt{-h} \left[\frac{h^{\alpha\beta}}{2\alpha'} (\partial_{\alpha} X) (\partial_{\beta} X) \right]$ $exp[-(x_1 - x_2)^2]$ X_1 X_3 Discretized World-Sheet

• Strings on Random lattice [Douglas,Shenker] $S = \sum_{ij} (X_i - X_j)^2$

 $\Rightarrow Z = \int \mathcal{D}h \ \mathcal{D}X \ e^{-S} = \sum \int d^D X \ \prod_{ij} e^{-\frac{1}{2\alpha'}(X_i - X_j)^2}$

Dual Field theory action

$$\hat{S} = \int d^D x \ tr \left[\frac{1}{2} \phi e^{-\alpha' \Box/2} \phi + G^{n-2} \phi^n \right]$$

Linear Regge trajectories: Confinement [Grisaru, Siegel, Y.T.]

Finite Order Gravity

Improved UV behaviour: 4th Order Gravity $S = \int d^4x \sqrt{-g}(R + c_0R^2 + b_0C^2)$ even Renormalizable [Stelle, 1978] Asymptotically free + Renormalizable!

Unfortunately $b_0 \neq 0 \Rightarrow$ Ghosts

If $b_0 = 0$ Asymptotic freedom, Renormalizability lost

Ghost + assymptotically) free gravity => NP gravity

Propagator

Scalar-Tensor Picture: HD terms in φ

$$S = \int d^4x \sqrt{-g} \left[e^{-\phi} R + \psi \sum_{0}^{\infty} c_i \Box^i \psi - \{ \psi (e^{-\phi} - 1) \} \right]$$

p-adic scalars in a curved background + dilaton?Field Equation

 $(1 - 6\sum_{0}^{\infty} c_i \Box^{i+1}) \phi \equiv \Gamma(\Box)\phi = 0 \Rightarrow \Delta(p^2) = \frac{1}{\Gamma(-p^2)}$

- Ghost free if Γ(□) has: a single zero, R² gravity no zeroes, Gaussian's
- Improved UV behaviour:

$$\Delta(p^2) = \frac{1}{(p^2 + m^2)}$$
$$\Delta(p^2) = e^{-p^2/m^2}$$
$$h \sim \frac{\operatorname{erf}(r)}{r}$$

Transition to FRW, $\Lambda = 0$

Late times

 $a(t) \rightarrow e^{\lambda t}$ & HD terms \rightarrow sech²(λt) ~ $e^{-2\lambda t} \rightarrow 0$

a(t)

=> Einstein Gravity & dS Universe => $\Lambda \neq 0$

Near Bounce

 $G_{00} \rightarrow 0$ but HD terms finite

Approximate Bounce

- Small times: HD terms = radiation
 We found ghost free examples
- Transition: HD terms ~ G₀₀
- Large times: FRW cosmology, HD terms << 0</p>

$$a(t) \sim t^{\frac{1}{2}}, \qquad G_{00} \sim \frac{1}{t^2}, \qquad \widetilde{G}_{00}^n \sim \frac{1}{t^{2(n+1)}}$$

Asymptotic Safety [Weinberg, 1976]

- Renormalizability replaced by asymptotic safety
- Quantum behavior captured by RG flow

$$\mu \frac{dg(\mu)}{d\mu} = ag^2(\mu) \qquad (a > 0) \qquad g(\kappa\mu) = \frac{g(\mu)}{1 - ag(\mu)\ln\kappa}$$

- Asymptotic safety = non-singularity \rightarrow UV fixed point 4 d gravity => G_N \rightarrow 0 asymptotic freedom
- Although ghost free finite HD gravity theories exist, (Ghost + asymptotically) free gravity => NP gravity
- Quantum Gravity actions (closed under renormalization flows) contains specific infinite series of HD terms: [Krasnov, gr-qc/0703002]
 Equivalent to "second order theory", no IVP

Hagedorn Phase

Qualitative Behaviour

- Close to T=T_H ~ M_S massive (winding) string states are excited. Pumping energy doesn't increase temperature, produces new states.
- Thermodynamics no longer determined by massless modes: $E = T_H S - bVT_H^{d+1} + ...$
- Cosmological Evolution: Entropy is constant
 Energy, Temperature remains approximately constant

$$\frac{\Delta T}{T_H} \equiv \frac{T_H - T}{T_H} \sim \exp\left[-\frac{E}{V^{2/3}T_H^3}\right]$$

• Transition to radiation occur when $S \sim VT_H^d$

Cosmological BCS Condensation

- BCS Theory
- Free theory, fermions filled upto Fermi-sea
- Attractive four-fermion coupling contributes to negative energy.
- Vacuum gets a negative shift with the formation of mass-gap. $L = -i\overline{\Psi}\gamma^{\mu}\partial_{\mu}\Psi -$
- Auxillary field
- Integrate the fermions, use mean field theory to get non-perturbative potential for Delta
- Gap Equation

Potential

• Trace equation (Lorentz gauge) $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$

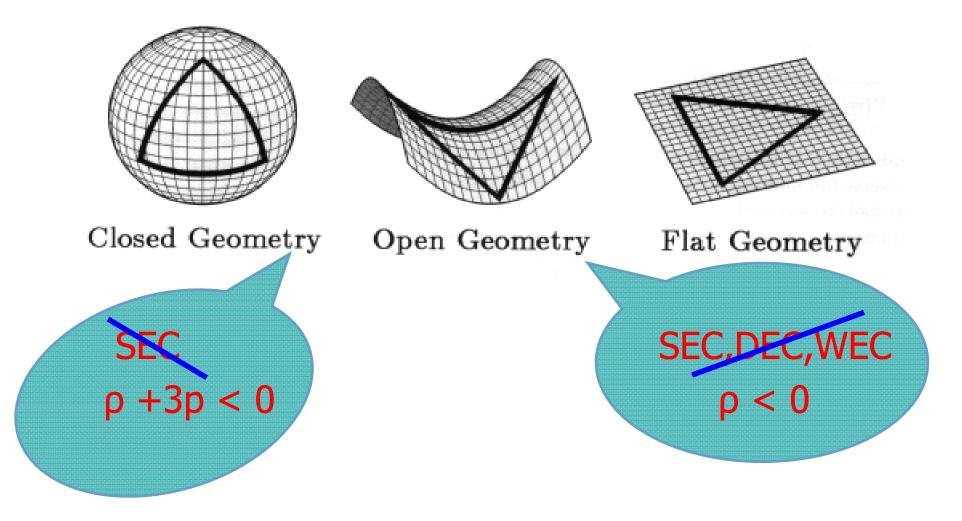
$$\Rightarrow \widetilde{G} = -\frac{1}{2} \Box (1 - 6 \sum_{0}^{\infty} c_i \Box^{i+1}) h = -\frac{1}{2} \Box \Gamma(\Box) h$$

Potential for h

$$\widetilde{G} \sim -m\delta(\vec{r}) \Rightarrow h(r) \sim \frac{1}{r} \int_{-\infty}^{\infty} dp \frac{p}{p^2 \Gamma(-p^2)} e^{ipr} \equiv \frac{G_N(r)}{r}$$

- AF: falls off exponentially $\Rightarrow G_N(r) \rightarrow r$
- Newtonian Limit $\Gamma(-p^2) \xrightarrow{p \to 0} 1 \implies G_N(r) \to const$.
- Example: $\Gamma(\Box) = e^{-\Box} \Rightarrow h(r) \sim \frac{\operatorname{erf}(r)}{r}$

Big Bang Singularity In GR at t = 0 we encounter a singularity R, $\Box R$, $\rho \rightarrow \infty$



<u>Non-singular Bounce</u>

- Ansatz: Find a(t) such that □R ~ R
 (...)R(t)+(...)R²(t) ~ matter sources
 Reduces to solving algebraic equation
- Hyperbolic Bounce a(t) = cosh(λt) works!
- Evolution

$$\widetilde{G}_{00} = T_{00} = \frac{1}{3} (\Lambda + \rho_{rad})$$

Solutions exist for asymptotically and ghost free theories: